MICHIGAN STATE UNIVERSITY

CYCLOTRON PROJECT

Cycletron Ion Starting Times In The Second Harmonic Accelerating Mode*

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ABSTRACT

CYCLOTRON ION STARTING TIMES IN THE SECOND HARMONIC ACCELERATING MODE

By

Herman Brenner White, Jr.

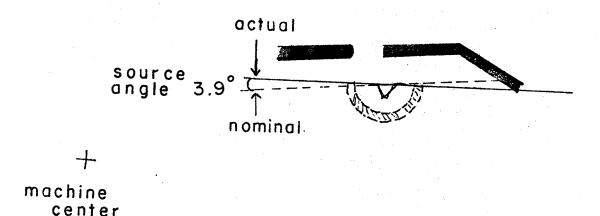
The analysis of orbit histories of the Michigan State
University Isochronous Cyclotron by computer calculations
is complicated in the central region because of the
strong electric forces. A study was begun to improve the
second harmonic operation utilizing the orbit code CYCLONE
and measured electric fields in the central region. Presented here is the experimental analysis of the orbit
calculations in terms of the correlation between predicted
beam positions by CYCLONE and experimentally determined
positions by foil burns. The data is presented graphically
to exhibit the close fit or deviation of the data with
various runs of the computer code, both for previously used
fields and for the new fields. Finally the work is extended
to help predict the optimum starting time and also therefore,
the starting phase of the particle beam.

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Note Added in Proof:

An angular alignment error between the puller gap and the ion source face has recently been observed; this error presumably existed at the time of the foil burn runs described herein. The ion source face was intended to be parallel to the face of the puller; telescopic measurements showed a deviation due to the calibration errors as shown in the sketch (viewed from the top):



The conclusions of this study are presumably slightly affected by this error; results should be interpreted accordingly.

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1. FORWARD

It has been often said lately, that most scientific work reflects very little of the human element behind it.

While the explanation of the work might very well be directed to the specialist in the field, the reasons for doing it and the research development behind it surely can be for the general reader. Luis Alvorez has pointed out that, "We should welcome. . .publications in which physicists tell us not only what they did, but why they did it, as well."

Therefore with this spirit and in this brief forward I should like to relate how this topic came about.

After having been involved with experimental modern (nuclear) physics as an undergraduate, I began work in trace elements here at Michigan State University using nuclear scattering techniques. This experience amply whetted my appetite for experimental work, and as a result I ventured off to other laboratories during the summer to increase my effectiveness in nuclear physics. While studying the $^{208}{\rm Pb}(\alpha, p)$ reaction at Argonne National Laboratory, I made frequent trips to the National Accelerator Laboratory at Batavia and it became apparent to me that a large part of our study of nuclear and particle physics could be found in the study of the accelerators themselves. One could say

understanding what the experiment does depends on understanding how the equipment works. In the following years I concentrated on accelerators and in particular the MSU Cyclotron. Prof. Blosser presented many questions about the machine (and answers I might add). The question of what happens to acceleration in the "Central Region" of the machine was one question some effort had been put into before. The central region being the area of the beginning history of the particle's acceleration. In this machine often times first harmonic operation has tended to be the Therefore, by studying this central region we felt that second harmonic acceleration could be improved. spending some time working on the CERN Synchro-Cyclotron, in magnetic measurements and particle detection, I began making electrical field measurements on the central region of the MSU Cyclotron. One important, and it turns out, very fundamental question arose: "How well can we predict, using these measurements, orbit properties of the accelerated beams in the Cyclotron?" Moreso, how accurate are our orbit calculations for second harmonic operation?

Preliminary to the final computations of the orbit properties, a study was necessary after the measurements to verify that the orbit calculations did indeed correspond to what actually happens inside the cyclotron machine. Therefore, it was necessary to correlate the calculated particle properties with observed results which is the intent of this study.

2. INTRODUCTION

The Michigan State University Sector Focused Cyclotron is a multi-particle, variable energy machine with a maximum proton energy of 56 MeV. The RF-system consist of two 138° dees and a 42° dummy dee which can be operated in either push-pull or push-push modes, over a frequency range of approximately 13.5 to 22 mega-cycles.²

The difficult analysis of orbit histories has in recent years been made easier by the use of the computer. The MSU Cyclotron Laboratory utilizes the Xerox Sigma-7 computer with which all the calculations for this study were made. The initial orbits calculated are crucial in determining the properties of the final extracted beam, such that the central region of the machine becomes of key importance. The complexity of the central region problem can be in part due to the interaction of particles with the magnetic forces, electrical forces, and space charge forces. Typically to study these forces one generally concentrates on accurately measuring the fields involved. Much attention is placed on the magnetic field outside of the central region, 3,4,5 since the particles are affected by magnetic forces here a longer period of time. That is, they are outside of the accelerating gaps longer being bent by the magnetic field

than inside the gaps being dominated by the electrical accelerating forces.

The electric field of the dees and the particular electric field configuration of the source-puller system becomes very important in understanding the initial orbit properties of the beam. Previous studies 6,7 have been done which discuss the central region geometry and orbit properties for previous electrical field measurements. Even though the theoretical calculations for particle trajectories and properties are based on physical laws, we still use these measurements as observed input data, to predict future behavior.

Since the last detail measurements of the electric field in the central region, some physical structure changes have occurred which could change the orbit properties of the cyclotron beam. To study these orbits it was necessary to make detailed measurements of the central region electric fields for the new geometry. Other measurements were made for the N=1 mode. Thus this study presents calculations for the N=2 mode to experimentally verify second harmonic orbit calculations.

3. PROCEDURE

3.1 Electrolytic Tank Measurement

The central region electric field measurements were done using an automatic electrolytic tank facility. The electric fields obtained were measurements of a rectangular grid of potential values in a two-dimensional array of points, ten square inches in size. The tank is characterized by a motor driven probe and the voltage of each point of the rectangular grid is automatically balanced to ground by a self-balancing bridge and punched on an IBM card. The scale is 1:1 for the actual machine.

The tank geometry is a scale model of the central region of the Cyclotron machine as shown in Figure 1. This geometry corresponds to the N=1 mode (push-pull) of acceleration with 134°-dees. The electrolyte used was doubly-deionized water, with the water-level in the tank at just above the top of the center electrode at the source position, to correspond to the median plane in the actual machine. A braking system was used on the tank's longitudinal movement to increase the amount of time for the bridge circuit to balance. This provides for a more accurate position of the probe at the time the voltage at that position is measured. Also oxidation and erosion of the copper surfaces

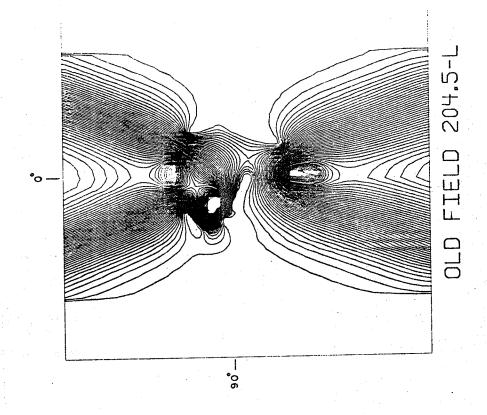
FIGURE 1.--Electrolytic tank model of the central region with the source and puller in the N=1 configuration. This structure shows one side of the median plane with the electrodes at the center defining the water level. Scale is 1:1.

were reduced by cleaning and the use of a dust cover. The cover also prevented conductive materials from contacting the water surface, as well as reducing motion of the water by air currents.

Another improvement was the larger size of the data grid, making possible more accurate calculations further out from the machine center. Consistant accuracy of ±0.5% was common in this measurement. All electric field data for orbit studies are stored in the laboratory computer library and coded for identification with a heading and a parameter list that gives all pertinent information concerning the potential grid that follows. The data (100x100 data points; large electric field) was first checked by plotting it with the computer in the form of an equipotential contour-map as shown in Figure 2. The utility of this procedure was to check that the geometry suggested was in fact reproduced in the gathered data. Also accuracy can be considered if all the equipotential contours are continuous and do not overlap, as they should provide analytical solutions to Laplace's equations for this complex geometry. equipotential contours are plotted for potential values representing 2% of the total range of potentials, that is from 2% to 98%. These contours can be compared with the N=2 contours from a previous measurement.

It should be noted that the electrolytic tank geometry in this field study is for the N=1 harmonic mode, but the

FIGURE 2.--Equipotential contour map for 134° dees from electrolytic tank data, in the N=2, push-push mode. The field coded 2.138.2A is the new field rotated by 45°. The old field, coded 204.5-L, contour map indicates the N=2 source to puller position with the puller recessed into the dee. In both cases these represent 2% contours of the total potential.



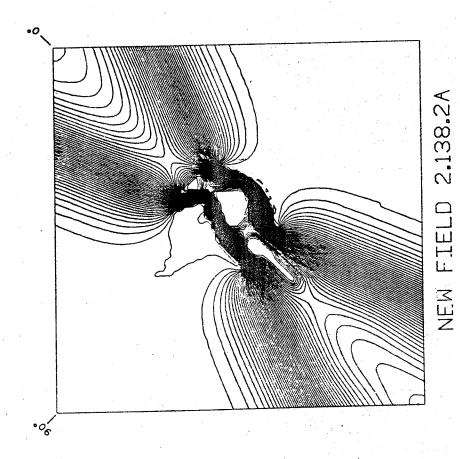


Figure 2

tank was run in the N=2 harmonic mode. The computer program used for these orbit calculations considers the forces due to the large electric field to not be included until the particle's angular position on the initial turn is greater than 90 degrees. The change in the position of the source and puller, however, is due to the phase dependence of the beam on the initial turns. That is, there is some starting phase (ϕ_0) for which the acceleration is a maximum, for which:

$$\int_{0}^{E_{\text{max}}} \sin \phi (E) dE = 0.$$

Since the magnetic field cone in the machine center shifts the particle's phase by approximately 21°, lengthening the first half turn compensates for this phase shift. 11

This puts the initial particle phase with respect to the rf near zero and results in the largest energy gain for the acceleration. Therefore, if the shift due to particles leaving the source on the peak of the rf cycle is approximately 21°, then recessing the source and puller into the dee by this amount will compensate for this phase shift. On the other hand, in the computer code the assumption is made that the particle only drifts after the source to puller field until it reaches 90° where the large electric field picks it up.

For this measurement the two dees have the same potential with one dummy dee grounded and the other having a potential 10% of the dees.

3.2 Source to Puller Field Measurements

The accuracy of the fields are important for accurate and reliable particle orbit calculations to be done later. For this reason special attention was placed on what is known as the source to puller region field. In this region the electric field is changing very rapidly and the accuracy of the electrolytic tank is not sufficient (with the total central geometry) to give very useful data in this critical region of the electrode geometry. This region characterizes the beginning motion of the ions to be accelerated. Since the magnetic field is comparatively very weak here, the electric field is the major focusing mechanism in the source to puller region.

The electrolytic tank was set to measure data at .1 inch resolution. Thus with the source to puller region of interest being approximately .50 square inches (on this scale), one would not get a reasonable resolution of the field in this area. To overcome this problem, a 20:1 scale drawing of the source and puller were made on conduction papers in much the same way as the 3 dimensional tank geometry. Electrodes were fixed to the paper by G. C. Electronics Silver conducting paint out-lining the perimeter of the electrode geometry. Using the Fluke 8200A Digital Voltmeter, a 15 inch by 20 inch grid with 0.025 inch spacing, of voltages was taken to constitute this smaller electric

field. This geometry's equipotential contour map is shown in Figure 3. For comparison the old electric field is also presented (1.06.7-S).

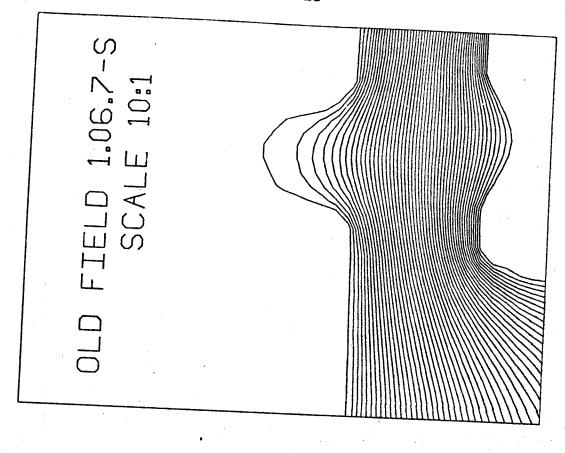
Definite improvements can be noted in these small field measurements. The larger scale made it possible to get more accurate data at the source slit. As can be seen, some physical changes were made from the previous 30° recessed source slit, mainly its position with respect to the puller gap. Previously, values at the source slit had to be interpolated to yield a continuous field because the number of data points in the slit itself was limited by the scale used. Therefore with a larger scale it was possible to obtain the necessary data points needed with a minimal use of interpolation.

3.3 General Outline of Foil Burning Technique

Predicted positions of the beam path can only be verified by physical measurements of the beam positions. To make these measurements the technique of foil-burning was employed.

A mesh of stainless steel, .0045 inches thick with 200 mesh squares per square inch, and 5.0 square inches in size was placed between the plane of the dees in the path of the accelerating beam. Foil holders were adapted to fit the current physical dimensions of the central region.

FIGURE 3.--Equipotential contour map for the source to puller electric field. Both 2% contours and the electric field configuration clearly shows the geometric change in the position of the source relative to the puller gap. Field 1.06.9-S was gathered on a 20:1 potential grid; the old field on a 10:1 grid. Both plotted 10:1 as shown, for N=2.



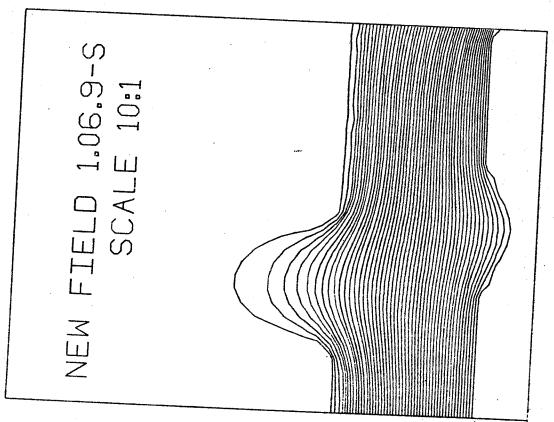


Figure 3

The foils were placed at 0°, 90°, and 270°. A differential probe trace was used in the 180° position, to define the beam path.

The machine was first set to yield a 24 MeV deuteron beam, and a complete differential probe trace turn pattern was done to get the total number of turns for the accelerated beam. This information is used to determine the approximate dee voltage. A vacuum break in the main chamber is necessary to place the foils, after which it is only necessary to restart the machine leaving the controls as they were, so that the parameters for the machine remain a constant. For this study centering coils were not used so that the magnetic field data would be in its simplest form for calculation.

As the beams impinges on the foil, the area of the beam heats the foil and eventually burns a hole at that radius, leaving the edges glowing as can be seen from Figures 4,5 and 6. The back of the copper foil holders acts as a beam stop so the particles don't accelerate further than the central region of interest which for this study was six turns. In measuring the 180° beam position, one limitation was power dissipation on the beam probe. After burning the foils we lowered the intensity of the beam to within a safe region. The power is given by:

 $P(watts) = E(MeV) \times i (\mu amps)/Z$

For our study at less than 10 inches radius with a beam current of 300 μA we dissipated approximately $7 \text{x} 10^{3}$ watts

FIGURE 4.--Foil holes near the 0° symmetry line of the machines. The glow comes from the beam burning the foils at 90° on the right and 270° on the left, for the 2nd run, of 24 MeV deuterons.

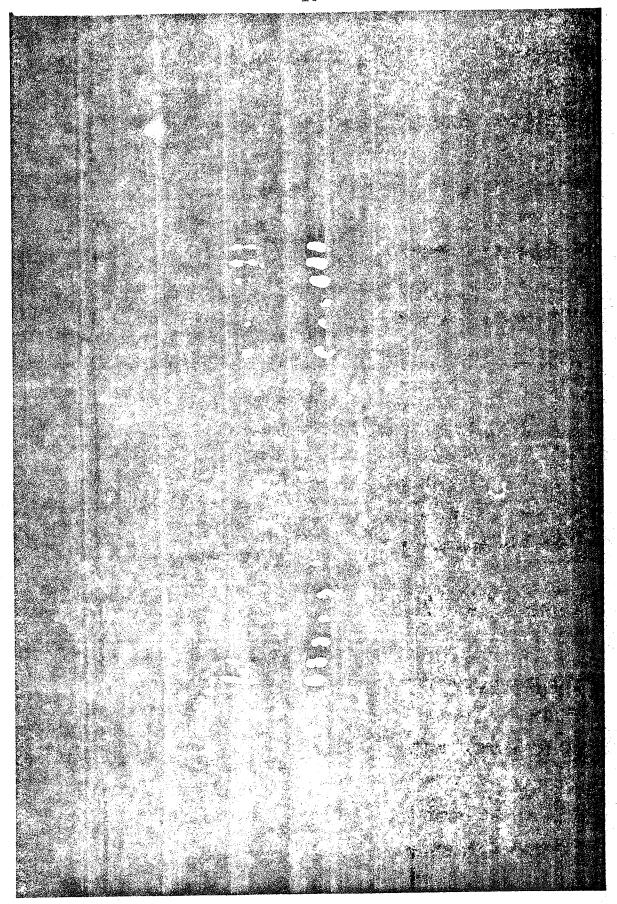


Figure 4

FIGURE 5.--Holes glowing along deflector angle. Position indicated by arrow in Figure 10. Photo shows first 5 turns at 0° and 4 turns at 270°, (1st turns at 270° blocked from view by foil holder at 0°).

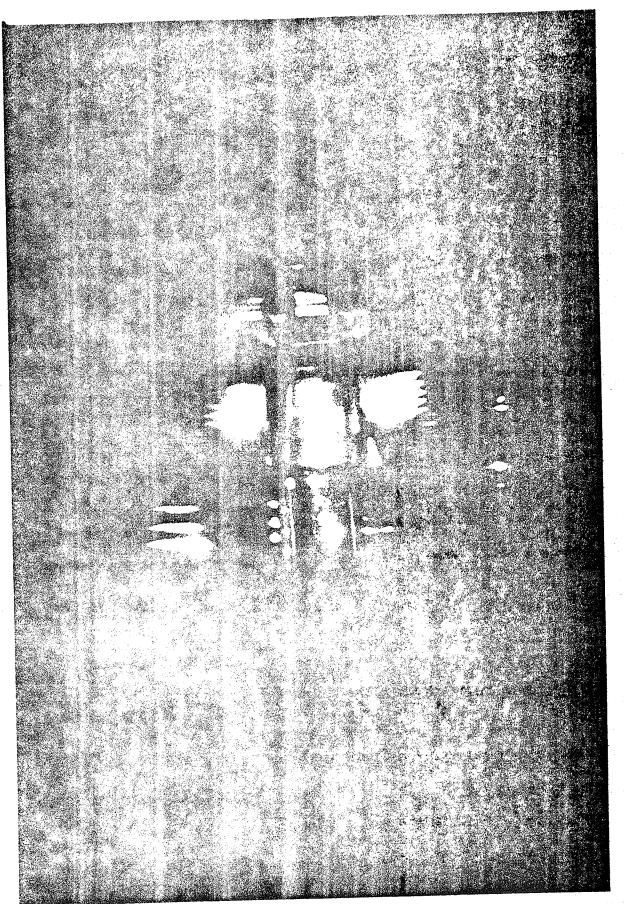


FIGURE 6.--Foil hole along detector angle as in Figure 5, showing the first turn on 0° just as the beam is about to burn through. Photo of 2nd run. Position of view indicated by arrow in Figure 10.

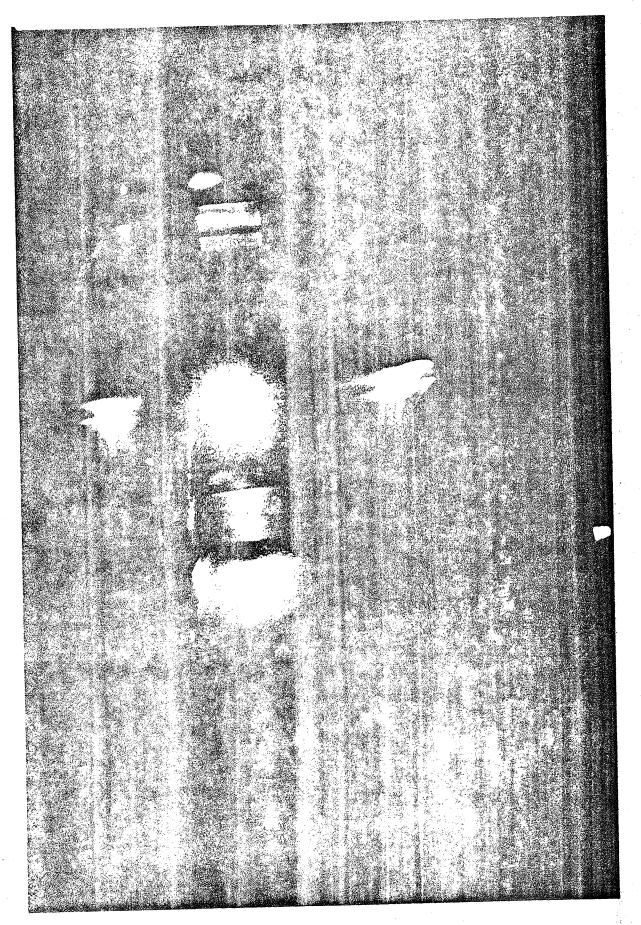


FIGURE 7.--Differential probe trace patterns. The probe is positioned at 180° and (a) shows the data at this angle giving the inner and outer turn radii. A full turn pattern with no centering coils in (b).

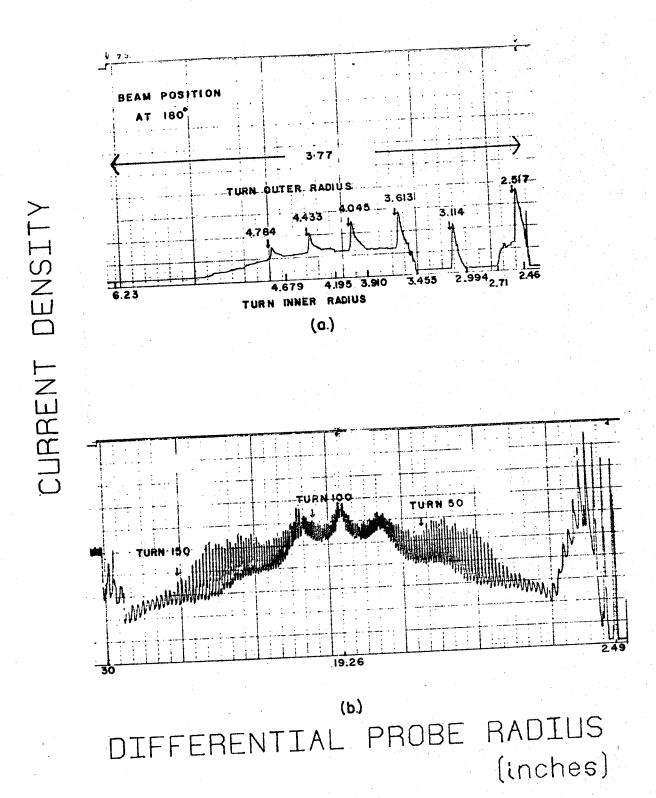


Figure 7

machine center was used as a radius vector. The position of each hole therefore was measured as a change in angle from the reference position (0°), with an assumed constant radius. Table II gives the radius of the holes, in the first run, with relation to the measured angles using the transit compass method. To check the data initially, the separation between each hole was compared to a vernier caliper measurement of the separation using the foil alone. Table II shows that this difference is as much as 0.1 inches. This method was abandoned because of these errors as well as other errors from slight motion of the transit. The data is also corrected on the south dee (90° position) for radial shift of the dees by .0625 inches due to the vacuum.

However, it was observed that each foil holder had been constructed with certain dimensions and alignment pins so they could be placed precisely, from the current cyclotron mechanical drawings. This method was pursued such that the position of the holders were known well with relation to the machine and therefore the 270°, 0°, and 90° data were measured with relation to the foil holders themselves. This arrangement is shown in Figure 8. The front edge of the copper holders was used as the reference point.

Table III and Table IV give the turn number, radius data, and angles for similar runs of 24 MeV deuterons, where the measurements were all done with respect to the foil holders. In Table IV, the 180° data is taken from the

TABLE II.--Foil hole radii at 90° and 270° for the first run determined by the transit compass measurements and calculations. Angle accuracy to 1 minute. The trigonometric radius vector used was 114.937 inches±.001 inch. South dee (90° data) corrected for radial displacement due to vacuum. Measurements relate to distance from machine center.

Position (degrees)		Radius	Calculated Separation	Measured Separation
		(inches)	(inches)	(inches)
	1°12'=1.2°	2.346	> .769	.623
	1°35'=1.583°	3.115	> .569	.524
	1°52'=1.867°	3.684	> .536	.415
90°	2°08'=2.133°	4.220	> .468	.371
	2°22'=2.367°	4.088	> .436	.355
	2°35'=2.583°	5.124		
	1°18'=1.3°	2.608	> .636	.734
270	1°37'=1.617°	3.244	> .502	.612
	1°52'=1.867°	3.746	> .435	.514
	2°5' =2.083°	4.181	> .368	.469
	2°16'=2.267°	4.549	> .469	.408
	2°30'=2.50°	5.018		

FIGURE 8.--The foil holders and their dimensions relative to the center of the Cyclotron at 0°, 90°, and 270°. The measurements are accurate to .001 inch. Note 90° and 270° are presented laying flat on the side viewed from the top.

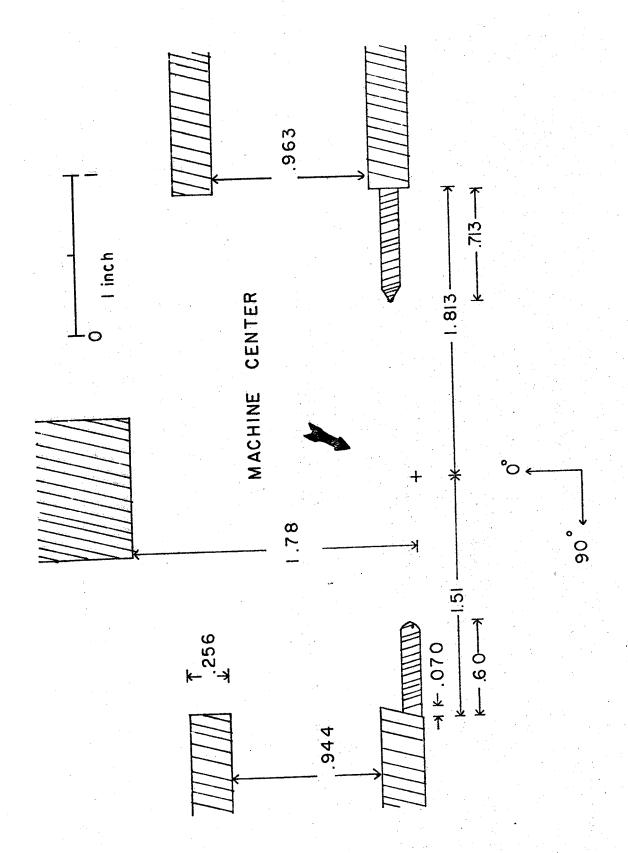


Figure 8

TABLE III. -- Hole measurements data for run 1 (24 MeV deuterons) obtained using the foil holders and mechanical drawing as references. Positions corresponds to geometry in Figure 8. Raw data from foil holders alone and final radii to an accuracy of .001 inch. No data shown at 180°.

		0°		90°	2	70°	
Turn #	Raw data	Radius	Raw data	Radius	Raw data	Radius	
0	-	-	_	_		_	
	-	_	-	·	.263	2.076	
1	.212	1.992	.432	1.942	.954	2.767	
	.312	2.092	.552	2.062	1.137	2.950	
2	1.029	2.809	1.183	2.693	1.576	3.389	
	1.127	2.907	1.316	2.826	1.766	3.579	
3	1.636	3.416	1.857	3.367	2.092	3.905	
	1.775	3.555	1.943	3.453	2.251	4.064	
4	2.163	3.943	2.362	3.872	2.530	4.343	
	2.291	4.071	2.462	3.972	2.693	4.506	
5	2.590	4.370	2.844	4.354	2.949	4.762	
	2.722	4.502	2.918	4.428	3.097	4.910	
6	3.016	4.796	3.265	4.775	3.354	5.167	
	3.112	4.892	3.337	4.847	3.461	5.274	
7	3.441	5.221					
	3.492	5.272	·		6.7		

TABLE IV.--Radii for run-2 using the foil holders and mechanical drawings as reference. Data at 0°, 90°, 180°, and 270° indicated. Raw data from foil holders alone and final radii all in inches (±.001 inch).

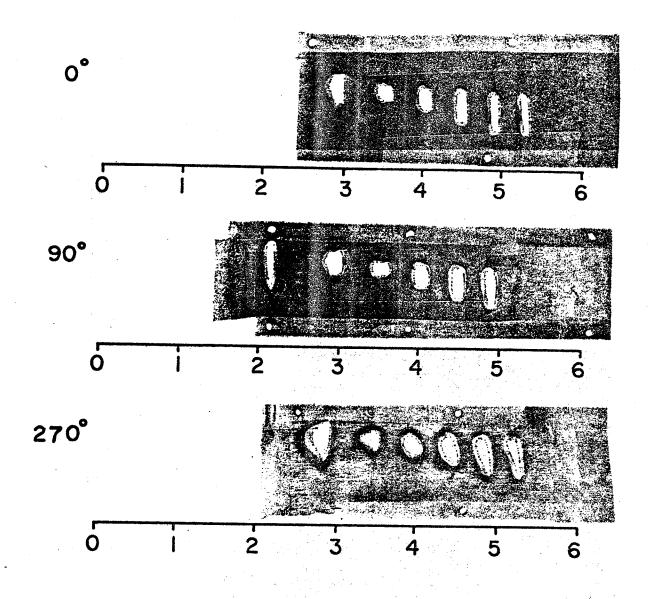
	0	0	9	0°	180°		70°
Turn #	Raw data	Radius	Raw data	Radius	Radius	Raw data	Radius
1		-	.572	2.012	2.517	.893	2.706
•		_	.650	2.090	2.710	1.087	2.900
2	1.057	2.837	1.303	2.743	2.994	1.497	3.310
~	1.234	3.014	1.521	2.961	3.114	1.710	3.523
3	1.668	3.448	1.906	3.346	3.455	2.031	3.844
	1.847	3.627	2.106	3.546	3.613	2.235	4.048
4	2.181	3.961	2.425	3.865	3.910	2.515	4.328
•	2.338	4.118	2.596	4.036	4.045	2.676	4.489
5	2.660	4.44	2.896	4.336	4.195	2.951	4.764
<i>y</i>	2.780	4.56	3.030	4.470	4.433	3.076	4.889
6	3.081	4.861	3.306	4.746	4.679	3.340	5.153
Ų .	3.173	4.953	3.438	4.878	4.784	3.452	5.265
7	3.465	5.245					
	3.558	5.338	e e				

differential probe trace pattern shown in Figure 7. differential probe gives the beam current at positions that are digitally displayed at the control consol. By knowing the beginning and end positions of the probe in the machine system one can measure the probe trace pattern, (measured as 8.29 inches) and taking the ratio of these numbers define a conversion factor that gives the position of all points on the probe trace pattern in units of the machine dimensions. The width of the peak was considered the radial length of the beam and the sharp edge of the peak is the outer radius of the beam. This procedure was verified visually by noting when the glow of the foil hole dissappeared when moving the differential probe from its outer most radius to inside the second turn. As was seen in Figures 4 and 9, some holes were flat on one edge and rounded on the other. Since the holes were not shaped the same, it was necessary to consider the width of the beam as the acceptable data.

Magnetic field data for input into the computer calculations was taken from the measured fields previously mentioned. A lagranian point-interpolation was done using the cyclotron settings to generate the raw field data, depending on the particle and its energy as well as the rf-frequency, main magnet and trim coils settings.

FIGURE 9.--Hole patterns burned on stainless steel screens, exposed at 0°, 90°, and 270° simultaneously.

On the 270° foil the screen was too short so the outside radius of the initial (0th) turn is slightly heated and burned, but not enough definition for data. (2nd run) 24 MeV deuterons.



RADIUS (inches)

Figure 9

4. DATA ANALYSIS

4.1 CYCLONE Computer Orbit Code

Beam orbit calculations for this study were done using the precise orbit code known as CYCLONE. This code utilizes exact median plane radial equations of motion for particle trajectories in crossed electric and magnetic fields.

Details of the electric fields are accounted for in three ways, which corresponds to the part of the program in which these various fields are used:

- (a) Source-puller region; in which the first part of the program considers the initial turn and uses the source to puller electric field.
- (b) "rf focusing" region; the second part of the program considers the first four turns and uses the large electric field;
- (c) the main acceleration region utilizes a stepfunction time-dependent potential in much the same way as the idea of assuming step function energy gain at each acceleration gap.

CYCLONE also provides provisions for study of equilibrium orbits, electrostatic deflector and magnetic channel. How-ever, none of these studies are included.

Out put from the CYCLONE program is provided in the appendices. Analysis of the CYCLONE program is presented here in the form of graphs, which compare the radius vs. turn numbers for various starting phases and radius vs turn number for the foil burned data presented earlier. In the use of this code it was necessary to pick the correct starting phase which would correspond to the foil burns. The method used initially was to run the program in succession changing the dee voltage until the beam is aligned with the radial slit on the initial turn at 1.713 inches radius for 180° as can be seen from Figure 10. However, the dee voltage can be closely approximated by using the differential probe trace to determine the number of turns and using the equation:

$$E_f/\eta = 4 V_{dee} \sin[N\theta_{dee}/2]$$

where E_f is the final particle energy.

η = number of turns

N = Harmonic number

 θ_{dee} = angular length of the dees.

Thus knowing the maximum energy gain per turn, it is possible to determine the approximate dee voltage.

FIGURE 10.--Plot of calculated orbit leaving the source at a starting time of -35° (rf-degrees), superimposed on a drawing of the N=2 central region geometry for 134° dees, in the MSU Cyclotron. Data experimentally gathered is indicated by the rectangles at all four positions, from foil burns and differential probe.

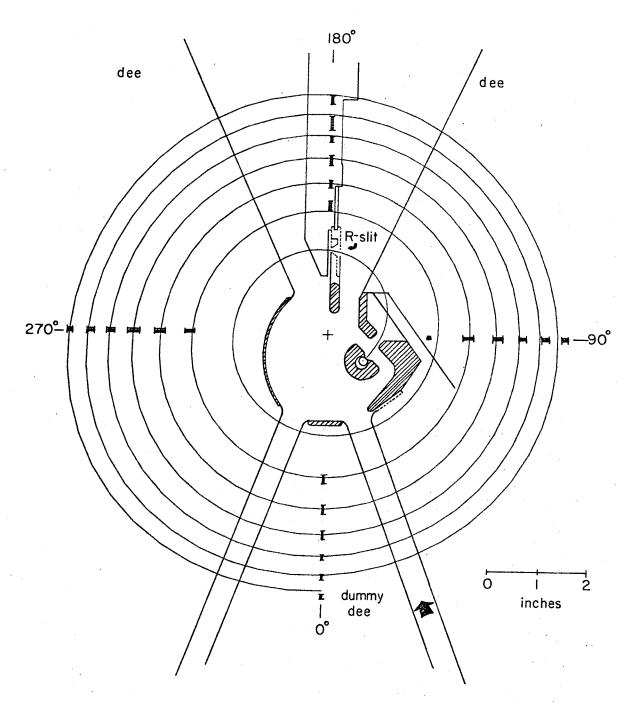


Figure 10

4.2 CYCLONE Calculations and Verification Analysis

In Figure 11, the results for one run are plotted against CYCLONE calculations for starting times (in rf-degrees) from -20° to -50° in 5° intervals. Figure 11a shows the results for the new fields and Figure 11b the results for the old fields. By comparing the width of the foil holes and the theoretical curve, one can compare which curves are accurate and therefore (since the only altered parameter are the electric fields), which fields correspond best to what actually exist in the Cyclotron machine. The width of the hole to some extent, is the degree of uncertainty in the beam position. Even though the measurements are accurate to within .001 inch, the center beam spot is still uncertain by the width of the burned hole. On the 0th turn only the outer edge of the hole is well defined.

Examining first the old field data (Figure 11b), it is found that at -20° starting time, a close fit is possible for data at 270° and 0° out to the fifth turn. In the sixth and seventh turns only the 0° data is closely fitted by the curve. This phenomenon is observed for the other starting times as well to such an extent that none can give a close correlation for all the holes at once. It becomes necessary to change starting times to fit the data only at a specific angle.

FIGURE 11.--Plot of inside and outside radii of foil data for the first data run, against turn number. Plot (a) gives theoretical curve from the CYCLONE code of radius-vs-turn number for various starting times using the new electric fields. The curves for the old fields are shown for comparison in (b). No data shown for 180°.

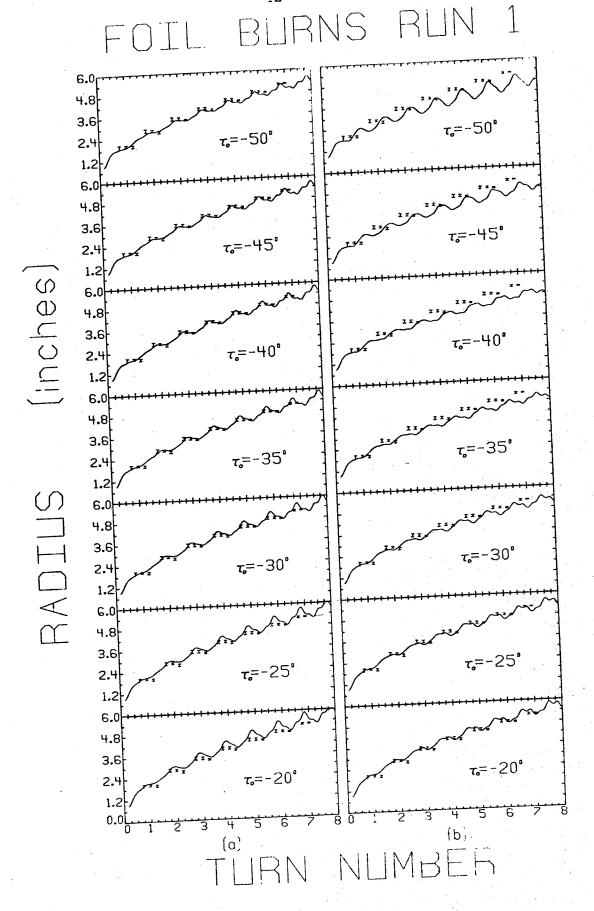


Figure 11

With the new fields (Figure 11a) the observation tends to be similar with the exception that for one starting time it is possible to fit accurately all the turns at all the data angles simultaneously, namely -35° and -40° . By plotting the different starting times for this run it is possible to observe the position of the Cyclone data shifting radially at the turns but still restricted to pass through the initial radial slit at 180°. The accuracy of the source to puller field and large electric field for the old data account for the shift in position according to turn number. In contrast Figure 12a fully illustrates the close correlation of Cyclone theoretical calculations to the measured field data. Data in Figure 12 also includes the 180° beam position making the analysis complete for four different sectors of the machine. Accurate agreement between experimental and theoretical positions was determined to be approximately .01 to .2 inches. Using the data from the second run this is illustrated better by introducing a parameter corresponding to the change in the position of the beam by Cyclone and the foil burns. fore we denote:

 $\Delta R = R_{\text{foil burns}} - R_{\text{cyclone}}$

as this change in position. Data for the foil burns, however, is presented as an inside and outside radius. Therefore, the centroid of this hole was considered the

FIGURE 12.--Plot of inside and outside radii of foil data for the second data run, against turn number. Theoretical curves for various starting times are plotted for the new fields (a) and the old fields (b).

FOIL BURNS RUN 2 6.0 4.8 3.6 2.4 τ=-50° τ=-50° 6.0 4.8 3.6 2.4 τ=-45° (inches) 1.2 τ_°=-45° 6.0 4.8 3.6 2.4 τ=-40° τ=-40° 1.2 4.8 3.6 2.4 τ=-35° τ=-35" 1.2 RADIUS 6.0 4.8 3.6 2.4 τ₀=-30° τ₀=-30° 1.2 4.8 3.6 2.4 τ₆=-25° τ.=-25" 1.2 4.8 3.6 2.4 τ₀=-20° τ₀=-20° 1.2 (a) THRI NIIMRFR

Figure 12

beam position and this was considered $R_{ extbf{f}}$ for the ΔR calculations. These calculations are presented in Table V thru Table VIII. In Figure 13 ΔR is plotted against starting time for successive turns, at each angular position of the By choosing a starting time at those turns where the line crosses 0.00 inches, it is easy to pick the initial starting time which gives close correlation to the foil data for this particular N=2 run. In this way a close determination of the initial phase at which the peak intensity portion of the beam leaves the ion source and consequently its phase dependence in the initial few turns can be studied. This is also a means of determining if the foil holes can be fitted at only one angular position at a time easily. In Figure 13 at 0° this seems true since the slope of the lines indicates a wide spread in starting time yielding a large range in data matching. The other angles indicate that -35° starting time should give accurate results for the second, third and fourth turns with a slightly larger spread at the fifth and sixth turns.

FIGURE 13.--AR-vs-starting time for the four burn positions, showing comparison of foil position to cyclone calculated radius. Each turn of the data is indicated. Note the greatly expanded radial scale contrasted to figures 11 and 12.

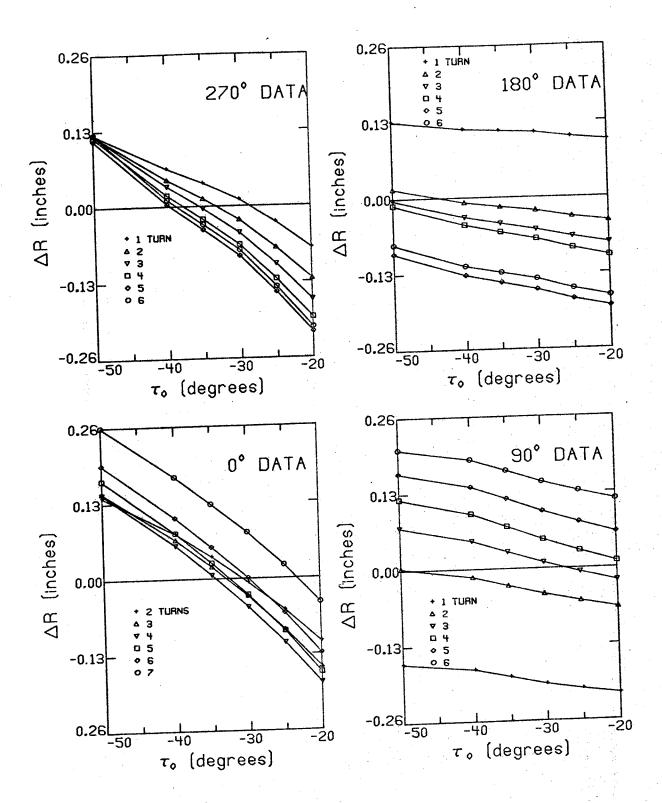


Figure 13

TABLE V.--ΔR-vs-turn number data at various starting times for the 0° data. Radii are presented in inches accurate to .001 inch.

τ_=-20°		τ _o =-25°		τ_=-30°	
Turn #	ΔR	Turn #	ΔR	Turn #	ΔR
2	10917	2	05425	2	00680
3	15193	3	08832	· 3	03048
4	17938	4	11047	4	04890
5	16167	5	09047	5	02808
6	12886	6	05992	6	.00073
7	04187	7	.02213	7	.07880

το	=-35°	τ _o =	40°	τ_=	-50°
2	.03717	2	.07729	2	.14044
3	.02008	3	.06690	3	.14726
4	.00598	4	.05696	4	.14535
5	.02654	5	.07793	5	.16879
6	.05357	6	.10421	6	.19491
7	.12753	7	.17430	7	.25954

TABLE VI.-- ΔR -vs-turn number data at various starting times for 90° data. ΔR are in inches $\pm .001$ inch.

$\tau_{o} = -20$	o° .	$\tau_{o} = -2$	5°	$\tau_{o} = -30$	
Turn #	ΔR	Turn #	ΔR	Turn #	ΔR
	21456	1	20647	1	19713
1 2	06761	2	05594	2	04354
3	02155	3	00672	3	.01021
4 .	.01094	4	.02952	4	.04973
5	.06107	5	.07913	5	.09795
6	.11762	6	.13233	6	.14858
					00
$\tau_{o}=-3$	35°	τ_=-	40°	τ ₀ =-5	U
	18418	1	17095	1	15976
1	02854	2	01317	2	.00300
2 3	.02899	3	.04830	3	.07280
4	.07222	4	.09483	4	.12131
5 5	.11978	5	.14101	5	.16504
6	.16822	6	.8740	6	.20568
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TABLE VII.--ΔR-vs-turn number data at various starting times for the 180° data. Radii are presented in inches accurate to .001 inch.

	=-20°	το	=-25°	τ_=	=-30°
furn #	ΔR	Turn #	ΔR	Turn #	ΔR
1	.09764	1	7007 =		ΔR
2	04084	2	.10317	1	.11120
3	07833	-	03301	2	02233
4	10045	3	06744	3	05468
.5		4	08683	4	07072
6	18726	5	17445	5	
O	16931	6	15635	6	15749
				O	13934

τ	o=-35°	1	-40°	т	=-500
1 2 3 4 5	.114490157404401058781449612820	1 2 3 4 5	.116990082103303045921325811696	1 2 3 4 5	.13138 .01655 00321 01125 09423 07866

TABLE VIII.--ΔR-vs-turn number data at various starting times for the 180° data. Radii are in inches ±.001 inch.

	τ _o =-25°	τ_=-	30°
$\tau_0 = -20^{\circ}$	Turn # ΔR	Turn #	ΔR
107183 212476 316068 419118 521657 620827	102733 207124 309939 412493 514746 613803	1 2 3 4 5	.011480232104513065770850407487
τ _o =-35°	τ _o =-40°		-50° .12410
1 .04004 2 .01288 300354 402175 503968	1 .06503 2 .04565 3 .03447 4 .01894 5 .00385	1 2 3 4 5 6	.12294 .12303 .11937 .11381 .12394

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5. CONCLUSIONS

It has been shown that this experimental procedure is useful in the prediction of a good starting time and thereby a good starting phase. The orbit radii and their correlattion to theoretical predictions seem to give a clue that a starting time of -35° to -40° gives a precise initial phase that best fits all the data. It is therefore concluded that this range for the starting time can be translated back to the actual machine thru the starting phase of the particles and thereby give better operation in the second harmonic mode.

This study of the central region electric fields has made it possible to substantiate our theoretical predictions. This investigation confirms that the new electric fields are accurate and reliable for future central region studies.

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APPENDIX

TABLE IX.--Cyclone calculations output for 24 MeV deuterons utilizing the new electric fields in the second run. Starting time is -35° (rf-degrees).

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TABLE X.--Cyclone calculations output for 24 MeV deuterons utilizing the new electric fields in the second run. Starting time is -40° (rf-degrees).

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