

# NEW EVIDENCE FOR PARITY INVERSION IN $^{10}\text{Li}$ FROM $^9\text{Li}-\gamma$ COINCIDENCES

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The structure of neutron-unbound  $^{10}\text{Li}$  remains an open question and thus also the complete understanding of the two-neutron halo nucleus  $^{11}\text{Li}$ . The knowledge of the  $(n+^9\text{Li})$  interaction is of essential importance to theoretical calculations of the three-body Borromean system  $^{11}\text{Li}$  ( $n+n+^9\text{Li}$ ) for which the two sub-systems  $^2\text{n}$  and  $^{10}\text{Li}$  are unbound.

Simple shell-models predict a  $\frac{1}{2}^-$  state as the ground state of  $^{10}\text{Li}$ , however the appearance of the  $\frac{1}{2}^+$  state from the sd-shell as the ground state of  $^{11}\text{Be}$  instead of the  $\frac{1}{2}^-$  was the first case of an intruder state in the shell structure of nuclei. This parity inversion has long ago been clearly observed experimentally where the non-normal parity state ( $\frac{1}{2}^+$ ) is 320 keV below the normal parity state which is the first excited state. Theoretical predictions [1] presented on Figure 1 show that the energy of the lowest non-normal parity state with respect to the lowest normal parity state increases systematically with Z for

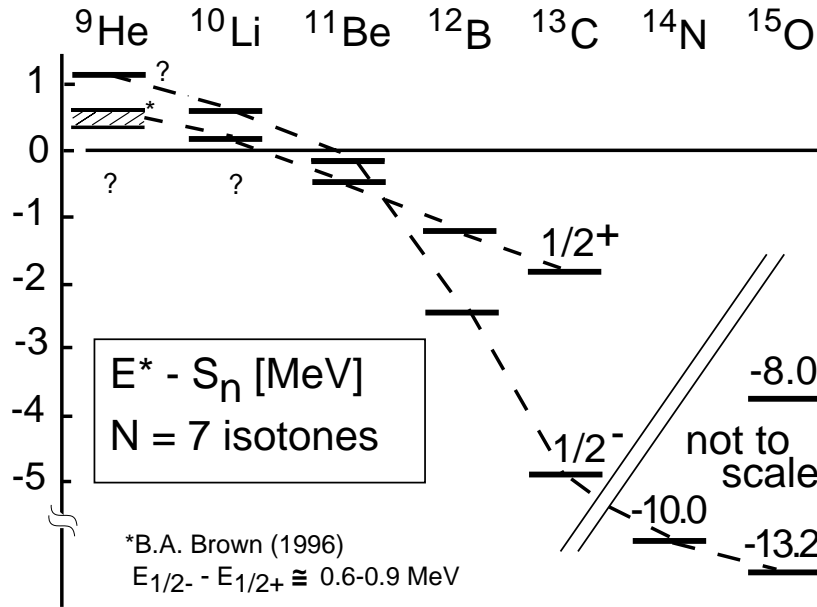


Figure 1: N = 7 isotone shell-model predictions for the lowest normal and non-normal parity states. Parity inversion is expected in neutron-unbound  $^{10}\text{Li}$  and  $^9\text{He}$ .

isotones with N = 7 and Z > 4. It is expected that this trend should continue in the other direction towards the lighter p-shell N = 7 isotones, and therefore the ground state of  $^{10}\text{Li}$  is predicted to be a neutron s-state, i.e. a non-normal parity state (several calculations [2-4] agree that this  $\frac{1}{2}^+$  neutron couples with the  $\frac{3}{2}^-$  proton to a  $2^-$  state).

Experimentally, the situation so far has been much more unclear. Several experiments using transfer reactions have shown evidence for a p-state around 240-540 keV, see Bohlen [5], Young [6], Caggiano [7] *et al.* The first experimental evidence for a neutron s-state at very low energy has been provided by the work of Kryger *et al* [8] using the method of sequential neutron decay spectroscopy at 0.

A virtual s-state at  $\sim 50$  keV was extracted corresponding to a scattering length of  $\sim -20$  fm. However the s-state could represent a decay to the first excited state in  ${}^9\text{Li}$  and would then not be the ground state. The observation by Zinser *et al* [9] of a narrow momentum distribution of the neutrons in coincidence with  ${}^9\text{Li}$  nuclei from the single neutron stripping of  ${}^{11}\text{Li}$  has been interpreted as final state interactions [10] and strongly suggested that the s-state couples to the ground state of  ${}^9\text{Li}$  (the removal of a halo neutron is not expected to excite the  ${}^9\text{Li}$  core).

An experiment was recently performed at the National Superconducting Cyclotron Laboratory of Michigan State University to investigate the possibility that the low decay energy s-state is not the ground state of  ${}^{10}\text{Li}$  and feeds into the first excited state of  ${}^9\text{Li}$  with the subsequent emission of a 2.7 MeV  $\gamma$ -ray.

A secondary beam of  ${}^{11}\text{Be}$  (46 MeV/nucleon) produced and separated in the A1200 spectrometer bombarded a  $300\text{ mg/cm}^2$   ${}^9\text{Be}$  target located in the N4 vault. A double-sided silicon strip detector located just after the target measured the energy loss ( $\Delta E$ ) and position of the charged fragments. Their residual energy ( $E$ ) was measured in 8 vertical plastic scintillator bars located at the exit of a deflecting magnet. Particle identification was achieved from the  $\Delta E$ - $E$  matrix, showing a clear  $Z$  separation. The mass number identification required a Monte Carlo simulation to determine the shape of the contours to apply onto the  $\Delta E$ - $E$  matrix. Neutron and  $\gamma$ -ray coincidences from the breakup of  ${}^{11}\text{Be}$  into the channels of  ${}^{10}\text{Li}$  and  ${}^{10}\text{Be}$  were measured using the MSU neutron walls at  $0^\circ$  and the  $\text{BaF}_2$  array of ORNL-NSCL-TAMU at  $90^\circ$ . The two liquid scintillator filled neutron walls covered an area of  $2 \times 2\text{ m}^2$  with a total efficiency of about 20 %. The 144  $\text{BaF}_2$  crystals covered a solid angle of approximately 1.5 sr with a total efficiency of about 5 % for  $\gamma$ -rays of 2.615 MeV ( ${}^{228}\text{Th}$ ).

The  $\gamma$ -ray and charged fragment coincidence analysis required an event-by-event reconstruction of the  $\gamma$  showers to add back the pair production escape peaks located in the nearest neighbouring crystals. Moreover, since the  $\gamma$ -rays are emitted by a fast moving source ( $\beta \approx 0.3$ ) Doppler shift correction was necessary.

The analysis of these coincidences was optimized for the  ${}^{10}\text{Be}$ - $\gamma$  data which could be easily interpreted using the known level scheme of  ${}^{10}\text{Be}$ . Clear evidence was obtained (see Figure 2-a). for a 3.37 MeV  $\gamma$ -ray from the  $2^+$  state and for a 2.90 MeV  $\gamma$ -ray from the decay of the  $2^-$  state at 6.26 MeV onto the first excited state ( $2^+$ ). A weaker indication of a  $\sim 6$  MeV  $\gamma$ -ray which had not been previously observed has been attributed to the decay of the  $1^-$  state at 5.96 MeV onto the ground state ( $0^+$ ). The study of the  ${}^{10}\text{Be}$ - $\gamma$  coincidences from the single-neutron stripping from  ${}^{11}\text{Be}$  should shed light on the wave function of the ground state of  ${}^{11}\text{Be}$ , the dominant part of which is known to be a  $1s_{\frac{1}{2}}$ . Sagawa *et al* [1] have suggested that the parity inversion in  ${}^{11}\text{Be}$  can only be explained if more complex contributions, such as the coupling of  $d_{\frac{5}{2}}$  to the  $2^+$  of  ${}^{10}\text{Be}$ , are considered. Our data, as well as those from M. Fauerbach *et al* [11], show that not only the core excitation to its first excited state should contribute, but also higher states such as the  $2^-$  and the  $1^-$ . Relative intensities for these  $\gamma$ -rays have been preliminary determined and their final interpretation for the structure of  ${}^{11}\text{Be}$  is on the way.

The number of  ${}^9\text{Li}$ - $\gamma$  coincidences (see Figure 2-b) correspond to a feeding of  ${}^9\text{Li}^*$  (2.7 MeV) in the single-proton stripping reaction from  ${}^{11}\text{Be}$  of  $4 \pm 1$  %, the dominant uncertainty coming from the poor fragment mass identification (a contamination of about 25 % from  ${}^8\text{Li}$  was determined with the Monte Carlo simulation mentioned above). This branching ratio compares very well with a recent shell-model calculation [12] which leads to the formation of  ${}^9\text{Li}^*$  in the single-proton stripping from  ${}^{10}\text{Be}$  of 5.2 %. As expected the number of 2.7 MeV  $\gamma$ -rays is small, therefore the triple  $\gamma$ -neutron-fragment coincidences did not provide sufficient statistics to be analyzed. Although, the angular neutron distributions showed,

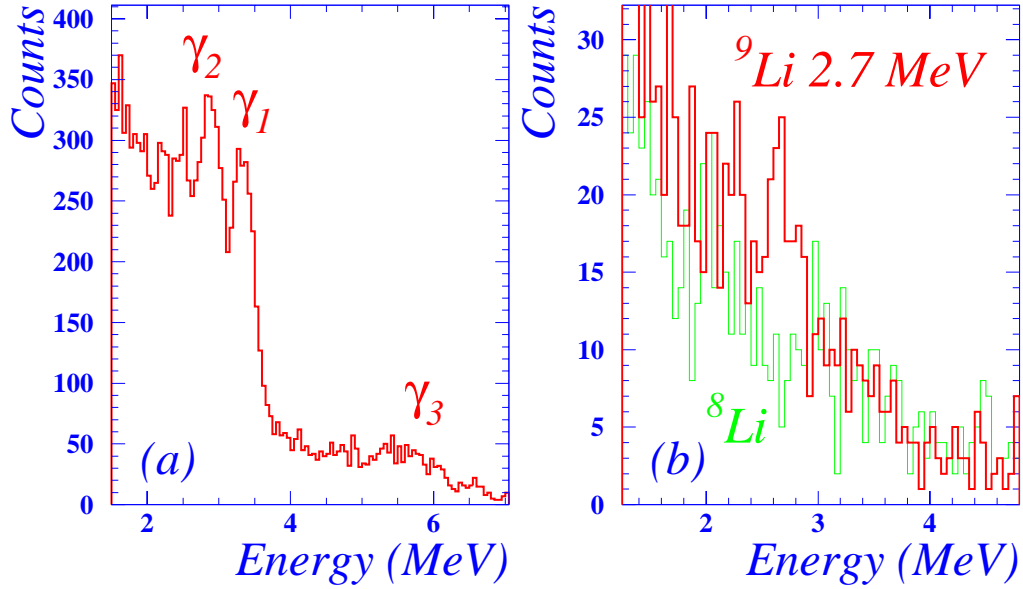


Figure 2: (a)  $^{10}\text{Be}$ - $\gamma$  coincidence spectrum showing evidence for 3.37 ( $\gamma_1$ ), 2.90 ( $\gamma_2$ ) and  $\sim 6$  MeV ( $\gamma_3$ )  $\gamma$ -rays. (b)  $^9\text{Li}$ - $\gamma$  coincidence spectrum (thick line) showing an excess of  $80 \pm 10$   $\gamma$ 's at 2.7 MeV in comparison to the  $^8\text{Li}$ - $\gamma$  coincidence spectrum (thin line).

similarly to the GSI and GANIL data, a dependence of the width with respect to the selected fragment: a narrower distribution in the case of  $^9\text{Li}$  resulting from final state interactions when a resonant intermediate state is formed in  $^{10}\text{Li}$  than in the case of  $^{10}\text{Be}$  corresponding to the diffraction dissociation of the neutron (this observation also supports the mass identification of the fragments).

Both, the weak population of the first excited state in  $^9\text{Li}$  and the narrow momentum distribution of the neutrons in coincidence with  $^9\text{Li}$  fragments indicating the formation of  $^{10}\text{Li}$  lead to the conclusion that the ground state of  $^{10}\text{Li}$  is an s-state which decays on the ground state of  $^9\text{Li}$ . Thus we observe, like in the case of  $^{11}\text{Be}$ , a parity inversion in neutron-unbound  $^{10}\text{Li}$ .

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#### References

1. H. Sagawa, B. A. Brown and H. Esbensen, Phys. Lett. B 309 (1993) 1.
2. B. A. Brown, Int. Symp. on Frontiers of Nucl. Struc. Phys., RIKEN, Japan (1993).
3. H. Kitagawa and H. Sagawa, Nucl. Phys. A 551 (1993) 16.
4. N. A. F. M. Popelier *et al*, Z. Phys. A 346 (1993) 11.
5. H. G. Bohlen *et al*, Nucl. Phys. A 616 (1997) 254c.
6. B. M. Young *et al*, Phys. Rev. C 49 (1994) 279.
7. J. A. Caggiano *et al*, to be published.
8. R. A. Kryger *et al*, Phys. Rev. C 47 (1993) R2439.
9. M. Zinser *et al*, Phys. Rev. Lett. 75 (1995) 1719.

10. F. Barranco, E. Vigezzi and R. A. Broglia, Phys. Lett. B 319 (1993) 387.
11. M. Fauerbach *et al*, in preparation.
12. N. Timofeyuk and I. Thompson, private communication.